The Crystal Structure of Aluminum Phosphate and Gallium Phosphate, Low-Cristobalite Type

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Crystals of the low-cristobalite phases of $AIPO₄$ and $GaPO₄$ have orthorhombic symmetry corresponding to the space group $C222$ ₁. The pseudo-tetragonal base-centered cell contains four formula weights of MPO₄. Parameters for each of the structures have been determined from the analysis of their X-ray powder diffraction patterns. Both structures are quite similar in arrangement to the analogous silica low cristobalite. Gallium is coordinated to four oxygens at a distance of 1-78 A. Aluminum is coordinated to four oxygens at a distance of 1.70 Å. The angle, M -O-P, is 135[°] in GaPO₄, and 145° in AlPO₄. The analogous angle in silica low cristobalite, Si-O-Si, is 150°.

Introduction

As part of a systematic investigation of the structures of the phosphates of trivalent cations, an X-ray diffraction analysis of galhum phosphate was undertaken. The compound was prepared by Alvin Perloff, who also studied the phase changes of the crystalline material between room temperature and 1100° C. The details of this work have been reported elsewhere (Perloff, 1956). The results relevant to the present discussion are briefly the following. The form of gallium phosphate which is stable at ordinary temperatures was found to be isostructural with $AIPO₄$, a well known low-quartz analogue. Two high-temperature modifications of $GaPO₄$ were observed. These are isostructural with two of the high-temperature phases of $AIPO₄$, which are identified as the low-cristobalite and the high-cristobalite types because of their close resemblance to the corresponding sihca modifications. These results are in agreement with the observations of other investigators (R. Roy, private communication).

This report is concerned with the structure determination of the low-cristobalite phases of $GaPO₄$ and $AIPO₄$, based on an analysis of powder diffraction data. Though the existence of low-cristobalite $AIPO₄$ has been known for a long time (Gruener, 1945; Beck, 1949), apparently no serious effort has been made to determine its atomic coordinates. This is understandable, since its powder pattern is almost identical to that of silica low-cristobalite, for which structural parameters have been found (Nieuwenkamp, 1935). Similarly, the pattern of low-cristobalite $GaPO₄$ is distinguished from the isostructural AlPO₄ pattern only by the effects of a small contraction in cell dimensions, and by such intensity differences as may be attributed to the presence of the heavier metal atom. Clearly, a model of low-cristobahte sihca on a gross scale should descriptively represent the structural arrangement of the other two crystals. Nevertheless,

the crystals cannot be identical in such interesting details as interatomic distances and bond angles. Though normally powder data would not be chosen for the elucidation of such details, the lack of single crystals of the compounds left no alternative. In this case, because of the simplicity of the structures, it seemed quite feasible to extract some dependable information from the powder patterns.

Experimental

Samples of the low-cristobalite phases of $AIPO₄$ and $GaPO₄$ were prepared in much the same way. Thoroughly ground samples of the low-quartz type were heated well above their high-cristobahte transition temperatures, and then allowed to air-cool quickly to room temperature. In both cases, the conversion to the low-cristobalite form is rapid and complete at the respective transition temperatures, and the phase persists unchanged at ordinary temperatures (Perloff, 1956). In order to avoid troublesome temperature effects, all of the diffraction patterns used for structure analysis were taken at room temperature. The material, which was fine-grained and well crystallized, yielded excellent powder patterns. All of the patterns were taken with filtered copper radiation by means of a wide-angle Geiger-counter spectrometer with a stripchart recorder. Sets of patterns were taken with scaling ratios of two, four and eight. The instrumental constants used were: receiving slit width, 0.006 in.; horizontal divergence, 1°; time constant, either 2 or 4 sec.; and scanning speeds, either $\frac{1}{4}^{\circ}$ or $\frac{1}{2}^{\circ}$ per minute; maximum counting rate, 180 per sec.

The patterns were indexed to a Bragg angle of about 75 ° , but the measurement of intensities was terminated at about 40° . At angles higher than this, excessive superposition of peaks made the measured intensities of doubtful value. In the useful range, numerical values of the relative intensities were obtained from measurement of peak areas, partially superposed peaks being resolved whenever feasible. Intercomparison of values from a number of diffraction patterns showed only minor variations, so the observations were presumed to be a reasonably reliable basis for structure determination.

The relative intensity values were corrected for the polarization-Lorentz factor, and, at a subsequent stage of the work, scaled to the calculated intensities by an appropriate constant. The temperature factor was negligible, as far as could be judged from the ratios of the observed and calculated intensities within the range of the observations.

Scattering factors were taken from the *International* Tables (1935) for Al⁺³, O^{-2} , P^{+5} , and slightly modified at low orders in the case of gallium to approximate the assumed state, Ga^{+3} . Atomic scattering factors might have been used without appreciable effects on the results.

Lattice type and space group

The unit cell of low-cristobalite silica, containing two formula weights of $SiSiO₄$, is primitive tetragonal, with dimensions $a = b = 4.96, c = 6.92$ Å (Nieuwenkamp, 1935). The pattern of $GaPO₄$, like that of $AIPO₄$, can be indexed to correspond to a closely analogous cell containing two formula weights of $MPO₄$. For AlPO₄, the dimensions are $a = b = 5.03$ Å, $c = 7.00 \text{ Å}, \text{ and for } \text{GaPO}_4, a = b = 4.93, c = 6.87 \text{ Å}.$ In both cases, the line positions give no evidence of deviation from tetragonal axial symmetry, though a small difference between the a and b axes quite possibly could be masked by the superposition of reflections at the higher angles.

The diffraction patterns of low-cristobalite $AIPO₄$ and $GaPO₄$ strongly resemble the pattern of the analogous silica crystal in intensity distribution as well as in spacing. There is, however, a difference in the type of extinctions which occur in the patterns. The space group of $SiO₂$ low cristobalite is $P_{4,1}^2 P_{1,2}^2 - D_4^4$ so that planes $00L$ are absent unless $L=4n$, and planes $H00$ are absent unless $H = 2n$. In AlPO₄ and GaPO₄, there is only one extinction; that is, $00L$ is absent unless $L = 2n$. These observations support the inference that the sites which are occupied in $SiO₂$ by an identical set of silicon atoms are shared, in the analogue structures, between phosphorus and metal atoms in regular alternation. Consequently, the $4₁$ and $2₁$ axes of the silica structure are transformed, respectively, to $2₁$ and 2 axes in the phosphate analogues so that only the systematic absence, $00L$, unless L is even, is required by their symmetry.

However, it is not necessary to predict a structure model, but only to assume the presence of discrete phosphate groups in order to deduce the space-group symmetry of the AlPO₄ and GaPO₄ low cristobalites. Only a few space groups are characterized by the single observed systematic absence, and simple considerations serve to eliminate all but two. In the tetragonal system, three space groups fulfill the required extinction condition; namely, P_{4_2} , P_{4_2} *m* and P_{4_2} 22. In each case, the metal and phosphorus atoms would have to occupy twofold special positions such that they would contribute intensity only to particular sets of reflections. The observed intensities completely refute this possibility. Therefore, the observed extinction must be due to a $2₁$ axis, and the crystal symmetry must be orthorhombic or lower, despite the tetragonal axial system.

Four space groups remain to be considered: in the orthorhombic system, $P222_1$ and $C222_1$ (with the necessary change in the choice of the unit cell); in the monoclinic system, $P2_1$ and $P2_1/m$. $P222_1$ can be eliminated on spatial considerations, and $P2_1/m$ on intensity considerations similar to those given for the tetragonal space groups. Only $C222₁$ and $P2₁$ remain as possibilities. Credible and, indeed, very similar arrangements may be based on either one. However, there is no reason to suspect that the coordination of the oxygens about the metal and phosphorus atoms is so irregular as to require the lower symmetry of $P2₁$ for its description; and subsequent detailed study of the intensities gave no indication of such a possibility. It appears, therefore, that the space-group symmetry of the low-cristobalite forms of $AIPO₄$ and $GaPO₄$ is $C222₁$, and that the true structural unit is not the small primitive cell heretofore used for comparison with that of low-cristobalite silica, but the larger basecentered cell, having as axes the diagonals of the original cell base, and containing four rather than two formula weights of $MPO₄$.

Crystallographic data

The crystallographic data for the low-cristobalite forms of $AIPO₄$ and $GaPO₄$ may now be summarized. The crystals are isomorphous. The base-centered cell, which is pseudo-tetragonal, contains four formula weights of $MPO₄$. The cell dimensions, determined from a small set of well resolved and moderatelyhigh-angle lines, are given in the following table, along with the calculated cell volumes and densities. The corresponding values for low-cristobalite silica, referred to analogous axes, are included for comparison.

It is interesting to notice that the cell volume of GaPO₄ is the smallest of the three, being 5.5% less than that of $AIPO_4$, and 2.2% less than that of $SiSiO₄$.

The space group $C222₁$ has eightfold general positions and two sets of fourfold special positions on twofold axes. The sixteen oxygens per cell may then be assigned to two sets of general positions, while the four Ga (or Al) and the four P atoms must lie on the twofold axes (in fourfold special positions). Following the *International Tables* (1935), the coordinates of the sets of equivalent positions are given by

$$
[0, 0, 0; \frac{1}{2}, \frac{1}{2}, 0] +
$$

-
- (a) 2 $x, 0, 0; \overline{x}, 0, \frac{1}{2}$ (4 P)

(b) 2 0, y, $\frac{1}{4}$; 0, \overline{y} , $\frac{3}{4}$ (4 Ga or Al) (b) 2 0, y, $\frac{1}{4}$; 0, \bar{y} , $\frac{3}{4}$
- *(c)* 1 $x, y, z; x, \bar{y}, \bar{z};$ $\bar{x}, \bar{y}, \frac{1}{2}+z$; $\bar{x}, y, \frac{1}{2}-z$ (80-1; 80-2).

Thus, eight parameters are required to determine the structure.

Details of structure analysis for low-cristobalite ALP04

The eight parameters required to fix the atomic positions must be deduced from powder diffraction data. The chief disadvantage of this type of data, namely, the superposition of non-equivalent intensities, is materially increased in this case because of the strictly pseudo-tetragonal nature of the cell; that is, pairs of reflections, *HKL* and *KHL,* are always coincident in the diffraction pattern. Further, all of the constituents of the cell have about the same scattering power, so that the sure location of one set of atoms independent of the others is not feasible. Therefore, the procedure must be one of trial and error; that is, the assumption of a complete structure, and the subsequent testing of its credibility by means of intensity calculations. The close relation between the powder patterns of the $AIPO₄$ and the low-cristobalite forms of silica certainly indicates that only small changes in the values of the $SiO₂$ parameters should be sufficient to convert them into reasonable values for the $AIPO₄$ analogue. However, the exact nature and extent of these changes are not too easily postulated. Almost any model based on acceptable interatomic distances gives some semblance of agreement between the observed and calculated intensities. Without a prior knowledge of the best agreement attainable from the available data, it seemed necessary to explore systematically, and in some detail, the entire range of parameter sets which could be considered reasonable.

The scheme adopted for this purpose involved certain assumptions: first, that the phosphate group is a regular tetrahedron, with phosphorus to oxygen distances of 1.56 Å; second, that oxygens of neighboring phosphate groups are at least 2-55 A apart; and finally, that the aluminum atom is equidistant from its four surrounding oxygens. On the basis of these assumptions, an extensive series of structure models was devised in the following way.

From the space-group symmetry it is known that phosphorus, the center of the phosphate tetrahedron, occupies a set of special positions $(x, 0, 0, \text{ and related})$ points) on twofold rotation axes of the crystal. **Conse-** quently, a twofold axis of the tetrahedron must coincide with the crystal axis, but the orientation of the group as a whole about this axis is unrestricted by symmetry. The phosphorus parameter, x , fixes the position of the center of the group on the axis; a second parameter, φ , may be used to specify its orientation. If φ is taken as the angle between any convenient reference plane—in this case, the cell base --and a given mirror plane of the tetrahedron, values of φ between 0 and 45° describe all structurally unique orientations of the tetrahedron. Thus, the required parameters are reduced to three, x for the phosphorus, φ for the oxygens, and y for the aluminum.

All values of φ greater than 25° may be eliminated from consideration at once, because they produce unreasonably small separations between the oxygens of different tetrahedra. For any given value of φ less than 25° , there is a small range of x values, increasing with decreasing φ , for which the O-O distances are acceptable. On the assumption that the aluminum is centrally located with reference to its four nearest oxygens, the aluminum parameter may be calculated for particular values of x in the permitted range. Thus, for any particular angular orientation of the phosphate group, it is possible to derive hypothetical structure models having consecutive small differences in the x parameters of the phosphate group, in the aluminum parameter, and consequently in the aluminum to oxygen distances. A tabulation of such parameter sets, for a series of φ values taken at suitably small intervals, gives a convenient means for studying the whole range of structural arrangements which are possible within the restrictions imposed by the preliminary assumptions.

Fortunately, a large number of cases can be eliminated by the use of a few simple planes as criteria. Some reflections can be found which are highly dependent on the orientation of the tetrahedron ; that is, on the y and z parameters of oxygen. For instance, the intensity calculated for 004 cannot be made to agree with the observed intensity unless the oxygen contribution is small and negative; the intensity of plane 222 depends completely on oxygen contributions, which must be fairly large. From these, and other less obvious examples, it appears that φ must be somewhere between 18° and 20° . For orientation parameters between 18° and 20° , the aluminum parameter may theoretically have values between 0.163 and 0.200 (unit-cell fractions). The corresponding Al-O distances range from 1.85 Å to 1.68 Å.

An attempt to use small groups of selected planes as criteria for narrowing the range of aluminum parameter values gave results which were not sufficiently decisive. Therefore, complete intensity calculations were made for model structures selected so as to sample the various combinations of parameter values as effectively as possible. A form of residual factor, $R_{\rm I} = \sum [I_o - I_c] \div \sum I_o$, was used as an estimate of the over-all agreement between the observed and eal-

culated intensities. This factor was found to change appreciably, and in a regular manner, for small differences in the aluminum parameter, and in the orientation parameter, as shown in Table 1.

The minimum value of R_1 attained was 11% ; attempts to reduce it still further by small changes, including slight deformations of the tetrahedron, were fruitless. The set of parameters for $AIPO₄$ which gives the most satisfactory agreement with the data is shown in Table 2.

Table 2. *Low-cristobalite parameters of AIPO*₄

A comparison of observed and calculated intensities is given in Table 3. The agreement is quite good. The largest discrepancies occur for planes of low intensity, where the measurement is most subject to error.

The calculated intensities of the pairs of planes, *HKL* and *KHL,* have been given separately in the table, though they cannot be so observed in the powder pattern. They have, in general, quite different intensities, so that single-crystal patterns of this crystal would be easily distinguishable by cursory inspection from single crystal patterns of low-cristobalite silica.

As has been shown, the parameters of low-cristobalite $AIPO₄$ were determined by a method which made no use of the known parameters of $SiO₂$. However, the set which gave the best agreement with the observed intensities is remarkably close to the set given by Nieuwenkamp (1935) for the silica crystal. His parameters, here referred to the base-centered rather than to the correct primitive cell for more direct comparison, are shown in Table 4.

The main difference between the structures is in the orientation of the oxygen tetrahedra; the parameter φ , which is 19° in low-cristobalite AlPO₄, is about 23° in the silica crystal.

Details of structure analysis of low-cristobalite GaPO4

The procedure used in the study of low-cristobalite $GaPO₄$ was essentially the same as that described for the analysis of the $AIPO₄$ crystal. The same preliminary assumptions were made, and a comprehensive series

Table 3. *Calculated and observed intensities,* AlPO₄ (low-cristobalite type)

$(\sin^2 \theta)_a$	HKL	$I_c/8$	$\boldsymbol{I_o}$
0.0236	110	$\overline{2}$	8
0.0357	111	2188	2020
0.0472	200 020	7 32 25	45
0.0485	002	7	11
0.0592	021 201	172 388 216	386
0.0721	112	689	621
0.0942	220	1080	
0.0953	202 022	89 1278 109	1391
0.1065	221	5	0
0.1170	310 130	$\bf{0}$ 2 $\overline{2}$	6
0.1301	311	152	
	131	189 354	411
0.1327	113	13	
0.1425	222	255 487	273
0.1562	203 023	794 307 ſ	760
0.1665	312 132	406 846 440	755
0.1885	400 040	46 104 58	114
0.1933	004	200	220
0.2002	401	13	
0.2030	041 223	12 602 I 577	540
$\boldsymbol{0\cdot 2125}$	330	0	0
0.2168	114	111	114
0.2246	331	879	
0.2264	133 313	1097 102 116	1193
0.2365	240	95	
	420 042	58 252 27	210
	402	72	
0.2401	024 204	$\boldsymbol{0}$ 23 23	20
0.2480	241 421	402 ١ 806 404 J	700
0.2606	332	1088	897
0.2843	242 422	338 589 251	700
0.2877	224	482	466
0.2978	043 403	258 508 250	553
0.3066	150	0	
0.3116	510 134	ı 954 370	1056
	314	583	
	151	260	
0.3190	511 333	209 677 208	706
$\boldsymbol{0}\text{-}\boldsymbol{3252}$	115	849	730

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	Table 3 (cont.)		
$(\sin^2 \theta)_o$	HKL	$I_c/8$	I _o
0.3447	243 423	516 738 1279	1030
0.3485	025 205	23 $\mathbf 2$	
0.3550	152 512	625 1167 542	982
0.3772	440	19	28
0.3819	044 404	155 285 130 í	363
0.3897	441	541	515
0.3960	225	80	81
0.4012	350 530	352 361 736	742
0.4055	334	23	
0.4128	351 531	108 152	
0.4152	153 513	332 299	
0.4194	135 315	424 402	
0.4244	060 600	2686 139 111	2476
0.4254	442	202	
0.4287	244	252	
	424	257	
0.4341	006	8	
0.4377	061 601	168 290 122	303
0.4489	352 532	18 122 104	151

Table 4. *Low-cristobalite parameters of* SiSiO4

of hypothetical structures were derived in the same way.

Because the cell of low-cristobalite $GaPO₄$ is somewhat smaller than that of its $AIPO₄$ counterpart, the maximum value of φ must also be smaller in order to avoid unreasonably short distances between neighboring tetrahedra. The limiting value is 15° in $GaPO₄$, compared to 25° in AlPO₄. The possible values of gallium parameters over the whole angular range of α fall between 0.150 and 0.200.

Because the scattering power of gallium is about twice as large as that of aluminum, there are no reflections from GaPO₄ which depend solely on the oxygen contributions, and few which are particularly sensitive to changes in the oxygen positions. Therefore, the orientation parameter cannot be limited to a small range of values at the outset, as was done in the case of low-cristobalite AlPO₄. On the other hand, many of the intensities are appreciably modified by small changes in the gallium position. By using groups of suitably chosen reflections as criteria, it is possible to show clearly that the gallium parameter must be somewhere between 0.180 and 0.190.

In order to refine the gallium parameter, and to find the associated oxygen positions, complete intensity calculations were made for models having a wide range of φ parameters, and a gallium parameter between 0.180 and 0.190. The value of R_1 changes in a small but consistent way for differences as small as 0-001 in the gallium parameter, with φ kept constant. It goes through a minimum for values of the gallium parameter between 0-180 and 0.185, attaining its lowest value close to 0.183. On the other hand, R_1 is not too greatly affected by the oxygen orientation. Table 5 shows three cases, calculated from the optimum

value of the gallium parameter and widely different orientations. Quite acceptable values of R_I are obtained for values of φ at opposite ends of its range. However, the value of $R_{\rm I} = 9\%$ seems to be a minimum, and the corresponding set of parameters, given in Table 6, were taken as the best obtained from the data.

Table 6. Low-cristobalite parameters of GaPO₄

	x	y	z
Gа		0.183	ł
P	0.327	Ω	o
	0.198	0.035	0.182
O_2	0.456	0.179	0.965

The parameters of $GaPO₄$ differ appreciably from those of $AIPO₄$ and $SiO₂$ (low-cristobalite forms); and the orientation angle of its oxygen tetrahedra is considerably smaller: 11° compared to 19° and 23° for the $A1PO₄$ and $SiO₂$ compounds, respectively.

Table 7 shows the calculated and observed intensities to be in satisfactory agreement.

Discussion of the structures

Owing to the obvious deficiencies of the data, the accuracy of the parameters cannot be high. However, the final sets of positions in each case give agreement with the observations which seems close to the limit set by the data, in the sense that a change in the gallium position by as little as ± 0.01 Å, and of the other atoms by less than ± 0.05 Å appreciably increases the factor, $R_{\rm I}$.

The interatomic distances and bond angles calculated from the parameters of the low-cristobalite phosphate structures are compared with those of loweristobalite silica in Table 8. The structure diagram

Table 8. *Distances and bond angles in MXO 4 (low-cristobalite type)*

of low-cristobalite $GaPO₄$ is shown in Fig. 1. With due allowance for such small differences as appear in the table, the same diagram could be used to illustrate the structural arrangement in $AIPO₄$ or SiO₂ (lowcristobalite forms). The fact that the three structures are quite similar can be inferred merely by inspection -6f their powder photographs. Therefore, the results of a structure analysis of $GaPO₄$ and $AlPO₄$ can have interest only in showing what differences exist between them.

In the low-cristobalite SiO_2 , the oxygens form a practically regular tetrahedron about each silicon. In the two $MPO₄$ structures, the phosphate group has been assumed to be a regular tetrahedron, and the distances, *M-O,* have been assumed to be exactly equal. The resulting configuration of four oxygens about M is a somewhat distorted tetrahedron, as can be seen from the six shortest 0-0 distances, which form the tetrahedral edges. This distortion may, of course, be exaggerated by the preliminary assumptions, if, as is quite probable, it is shared between the gallium-oxygen and the phosphorus-oxygen configuration.

Fig. 1. Gallium phosphate, low-cristobalite type; [001] projection of the crystal structure. The true base-centered cell is indicated by full lines, and the pseudo-cell, comparable to that of low-eristobalite silica, by heavy broken lines.

In low-cristobalite $SiO₂$ the angle Si-O-Si is 150°. The analogous angle is smaller in the two phosphates, and decreases with the heavier metal: it is 145° for Al-O-P, and 135° for Ga-O-P. This difference in the orientation of the oxygen tetrahedra accounts for the facts that the $AIPO₄$ cell is not quite as large compared to that of $SiO₂$ as might be expected from spatial considerations; and that the structural unit of $GaPO₄$ can be accommodated in a smaller cell volume than that of AlPO₄ or even SiO₂, despite its larger M -O separation.

No preliminary assumptions were made concerning the magnitude of the \overline{M} -O distances in the \overline{M} PO_c structures; a wide range of values were tested with various phosphate group positions. In the case of lowcristobalite \widehat{A} lPO₄, the distance Al–O was found to be 1.70 Å. This value has also been reported for fourcoordinated Al-O in the crystal of $AIAsO₄$ (Machatski, 1935), which is structurally analogous to low-quartz. It is considerably smaller than would be deduced from the assumption of ionic radii. For instance, in Al_2O_3 (corundum) the average O-A1 distance in six-coordination is 1.91 Å (three at 1.89 , three at 1.93 Å). Correction of this distance to four coordination gives about $1.79 \; \text{\AA}.$

The value of 1.78 Å was found for the Ga-O distance in low-cristobalite $GaPO₄$. There are no available experimental data with which to compare it. For ionic radii, the distance for four coordination would be in the neighborhood of 1.85 A.

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A Note on the Debye-Waller Theory

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The effect of temperature on the Bragg reflexion from crystals is examined theoretically for a general crystal. It is shown that measurements of this effect can be used to obtain information about the vibrational spectrum of a crystal. A comparison is made of the existing data and the theoretical predictions.

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The effect of temperature on the reflexion of X-rays from crystals was first investigated by Debye (1914), and in the form due to Waller (1925) is part of standaxd X-ray theory. The theory shows that the intensity of X-rays reflected by a particular set of Bragg planes varies with temperature in the form $\exp(-2M)$, where

$$
M = 8\pi^2 u_z^2 \sin^2 \theta / \lambda^2 \,, \tag{1}
$$

 θ being the Bragg angle, λ the wavelength of the X-rays and $u_z²$ an averaged mean square displacement in the direction z perpendicular to the planes. The